

Summary on transport code descriptions

→ Remarks on the nature of discrepancies between transport codes

physical input
(EOS, σ_{inmed} ,
 $\pi\Delta$ physics, ..)



**transport
code**



observables

- unique?, e.g. like 2N transfer
- very complex, simulation of an equation rather than a solution
- depends on the question you ask

Transport theory based on a chain of approximations

Martin-Schwinger hierarchy in many body densities:
truncation, introduction of self energies (1-body quantities)

Quantum transport theory: Irreversibility, Kadanoff Baym theory

semiclassical approximation :

Wigner transform, not necc. Phase space probabilities

Gradient approximation (sep.of short and long scales)

Quasiparticle approximation

Spectral function \rightarrow delta function with effective quantities

\rightarrow BUU equation

$$\frac{\partial f}{\partial t} + \frac{\vec{p}}{m} \vec{\nabla}^{(r)} f - \vec{\nabla} U(r) \vec{\nabla}^{(p)} f(\vec{r}, \vec{p}; t) = \int d\vec{v}_2 d\vec{v}_1 d\vec{v}_2' v_{21} \sigma_{12}(\Omega) (2\pi)^3 \delta(\vec{p}_1 + \vec{p}_2 - \vec{p}_1' - \vec{p}_2') \\ [f_1' f_2' (1 - f_1)(1 - f_2) - f_1 f_2 (1 - f_1')(1 - f_2')] + \delta f(r, p, t)$$

6-dim integro-differential equation, non-linear

\rightarrow simulate solutions

introduces many technical details

fluctuations
variance of 2b collisions
neglct of higher orders

methods of solutions:

$$\frac{\partial f}{\partial t} + \frac{\vec{p}}{m} \vec{\nabla}^{(r)} f - \vec{\nabla} U(r) \vec{\nabla}^{(p)} f(\vec{r}, \vec{p}; t) = \int d\vec{v}_2 d\vec{v}_1 d\vec{v}_2' v_{21} \sigma_{12}(\Omega) (2\pi)^3 \delta(\mathbf{p}_1 + \mathbf{p}_2 - \mathbf{p}_1' - \mathbf{p}_2') \\ [f_1' f_2' (1-f_1)(1-f_2) - f_1 f_2 (1-f_1')(1-f_2')]]$$

Boltzmann-Vlasov-like (BUU)
solve as exactly as possible:
- test particle method
 exact in the limit of $N_{TP} \rightarrow \infty$
- deterministic, no fluctuations
 include fluctuations explicitly
- connection between U and σ
 by approx of self energy,
 e.g. Brueckner theory

Molecular dynamic-like (QMD)
- inject classical fluctuations
 and correlations (nucleon wave packet)
- damped (finite Gaussians,
 averaging width Δx , parameter
+ Pauli correlations (AMD)
- relation between U and σ not so clear,

biggest difference:
role of fluctuations
fragmentation, correlation functions
but also affects Pauli blocking and collective excitations

Fluctuations: almost a „fight“ between MD and Boltzmann models:

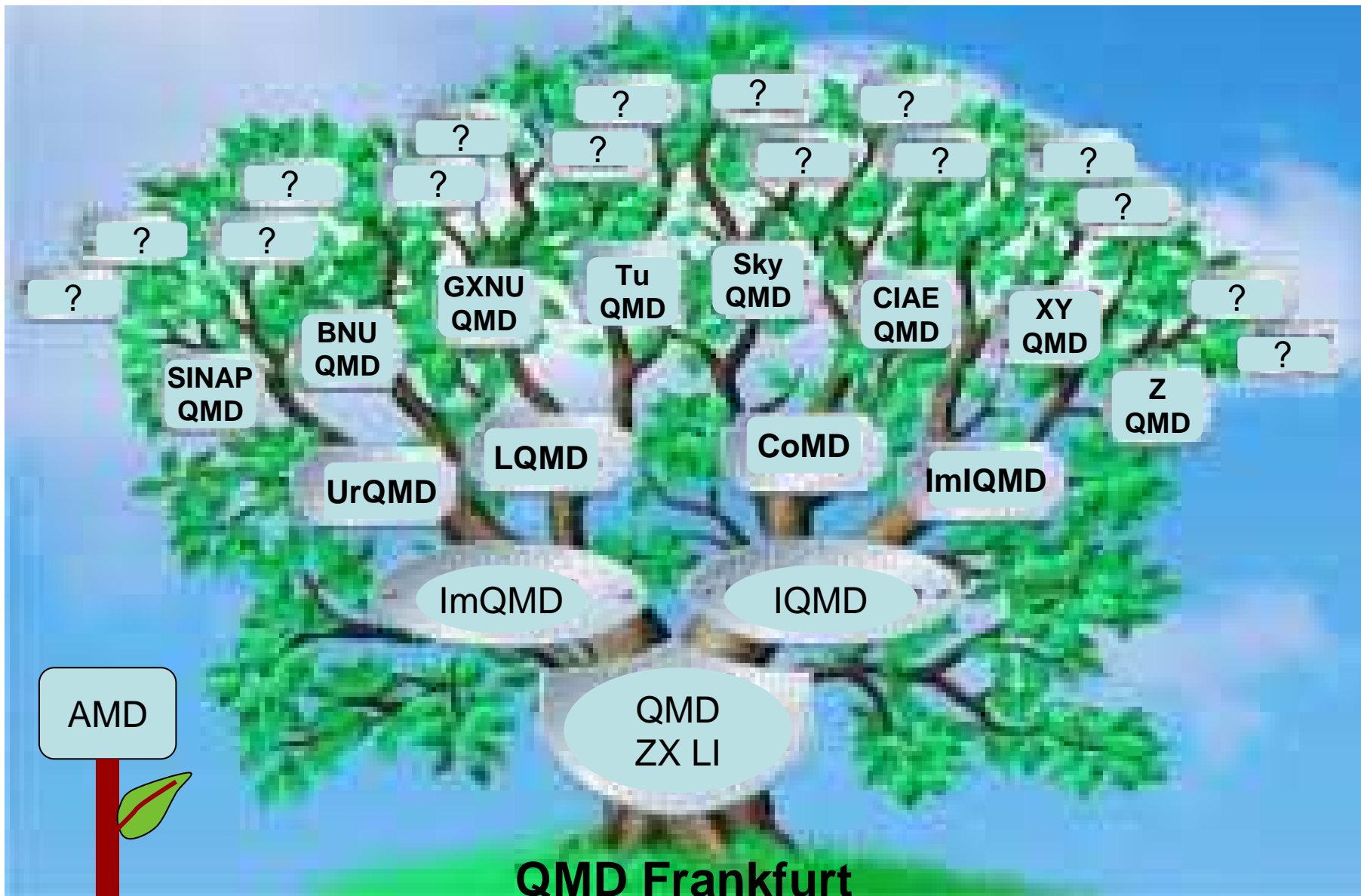
QMD!!!

$$\frac{df}{dt} = I_{coll} + I_{fluc}$$

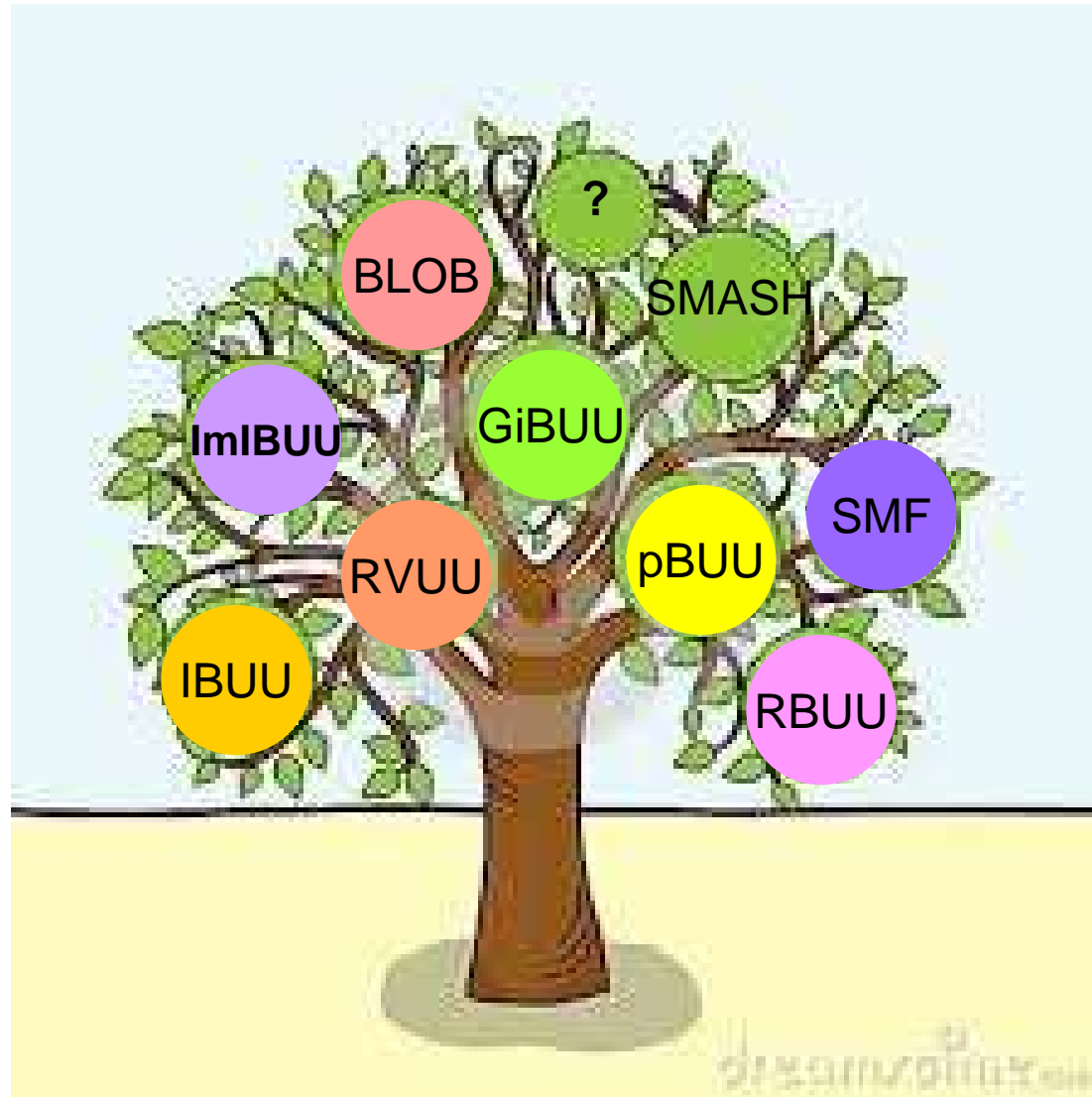
BUU??



now discussed beyond ideological barriers



„... in full bloom...“ – a good sign for the expanding activity,
 but try to make realtion and changes transparent,



„...lots of individuals...“

Steps in solving transport simulation

- initialization
- propagation of (test) particles (Vlasov)
- Collision partners and probabilities, elastic (Boltzmann)
- Pauli blocking (Ühling-Uhlenbeck)
- inelastic collisions (new particles), often perturbative, dep. on energy

Code comparison:

- differences of results of codes, e.g. isospin diffusion, pion ratios
- 1. phase: comparison of HIC with controlled input
 - differences seen (talk of Betty)
 - indications of reasons (initialization, Pauli blocking)
 - but difficult to pin point
 - general systematic theoretical error (30% (100 MeV), 13% (400 MeV))
how to improve?
- 2. phase: box calculations
 - better controlled conditions
 - exact limits often available
 - resolve differences because of strategies and of errors
from intrinsic differences (like BUU vs. QMD)

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initialization: solvable,

- initialize consistent with density functional used in transport
so that initial nucleus is a good approximation to the ground state
- more important than having identical density distributions

propagation: hamiltonian eom, easy

but

fluctuation dampen critically collective motions
momentum dependence, energy conservation

Time evolution of Fourier transform ρ_k

Second formulation of Homework #2:

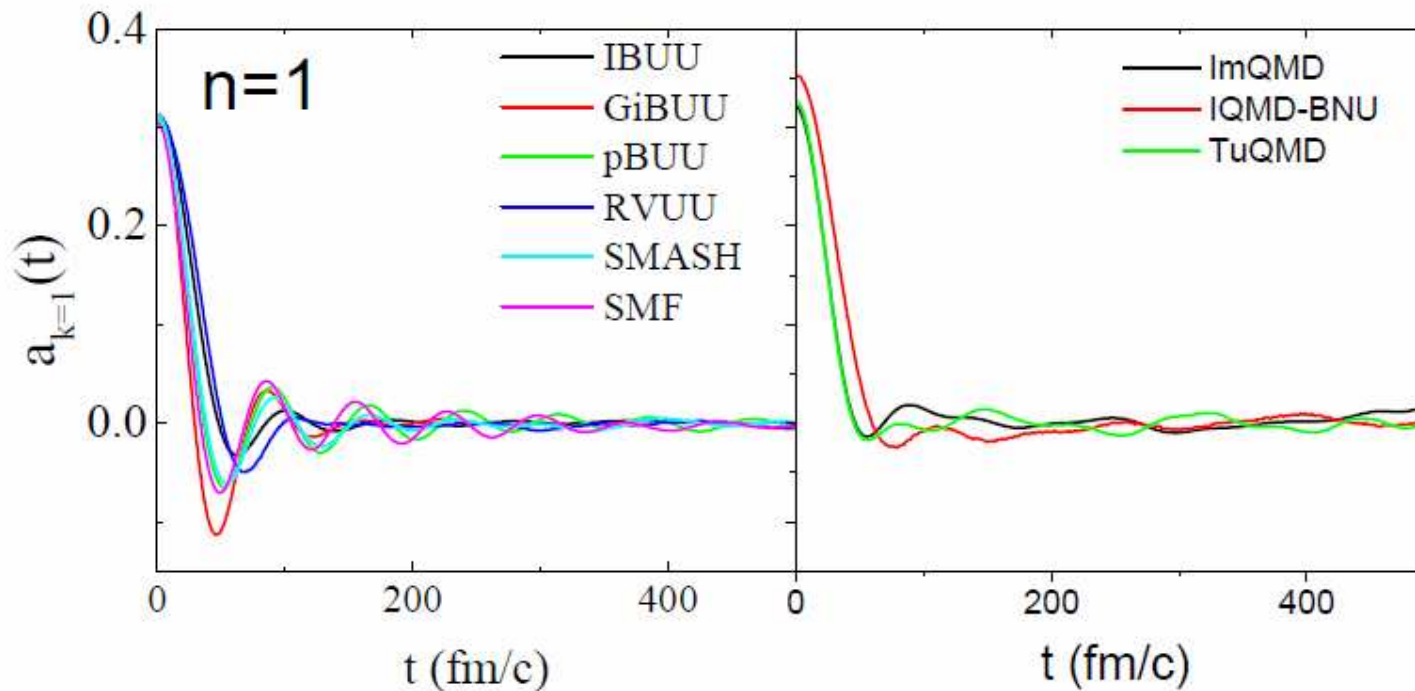
Longer final time and results given each 0.5-1 fm/c

n = 1

$$\rho_k(t) = \int dz \sin(kz) \rho(z,t) \quad k = n 2\pi/L$$

*Larger damping
and structureless fluctuations
In QMD-like*

Different oscillation frequency in BUU-like



Collision probabilities:

Bertsch prescription: particles collide,

- if their distance is below the interaction length and
- if they reach the distance of closest approach in their time step
- improve: the same nucleons should not collide again in the next time step

lesson: exact results come from kinetic theory, which makes assumption

in complete independence of collisions and equilibrium

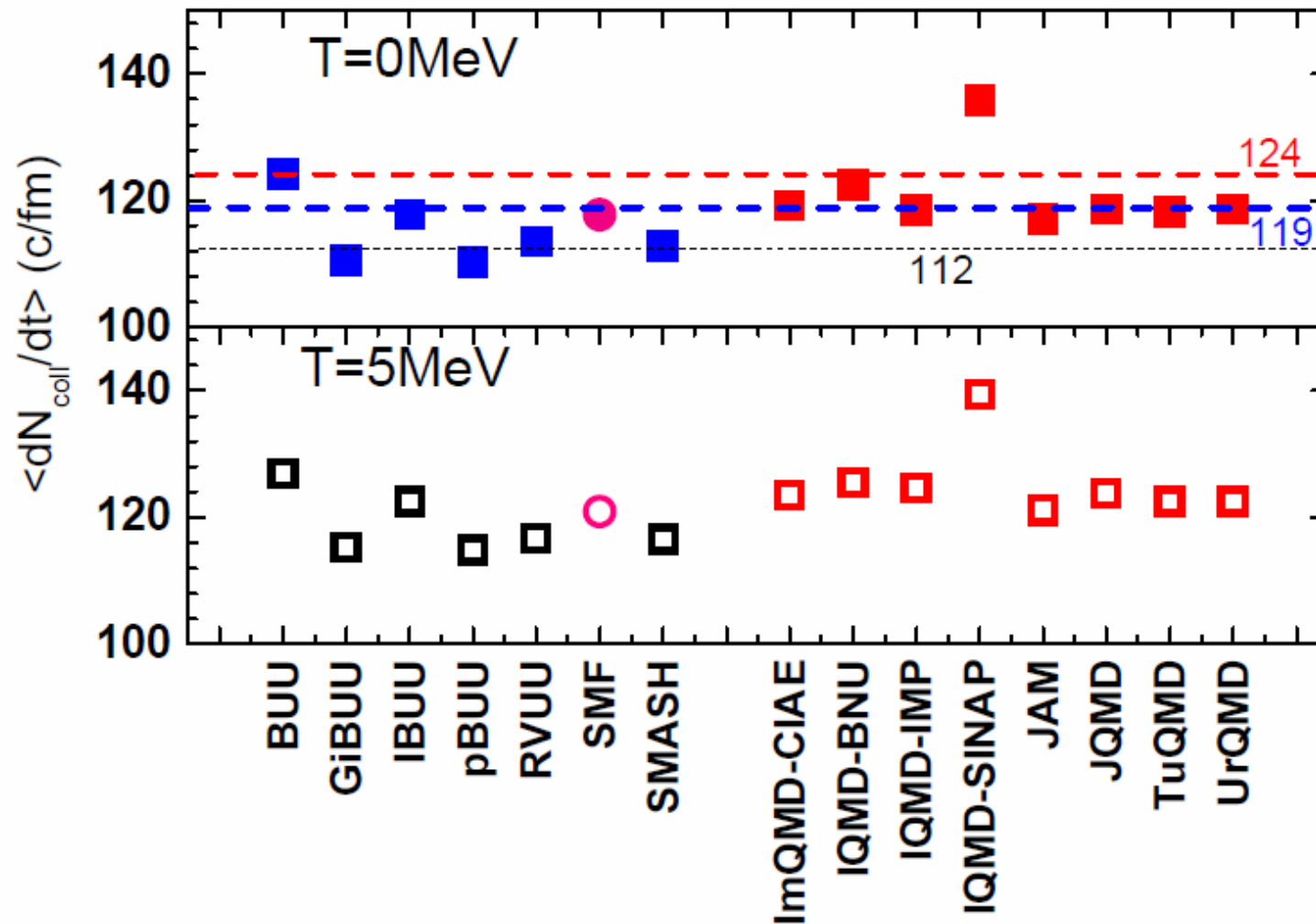
→ not so easy to follow in simulations (not always good)

mean free path description: assure mf path from kinetic theory

assure agreement limits put perhaps oversimplified in collisions (no equilibrium)

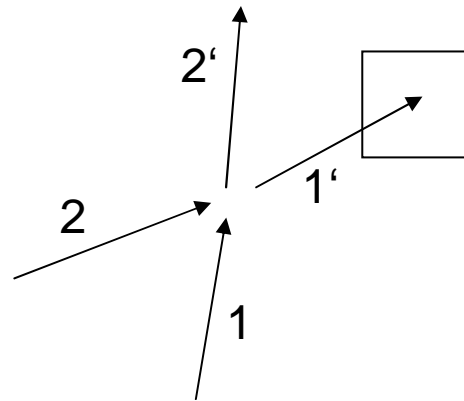
Theoretical results for CT0

Effect in dN_{coll}/dt for CT0: 124¹ (nonrelativistic) \rightarrow 119 (relativistic)



Ideal value for Boltzmann: 116.8 (nonrelativistic) \rightarrow 112.6 (relativistic, by J. Xu)

Pauli blocking:



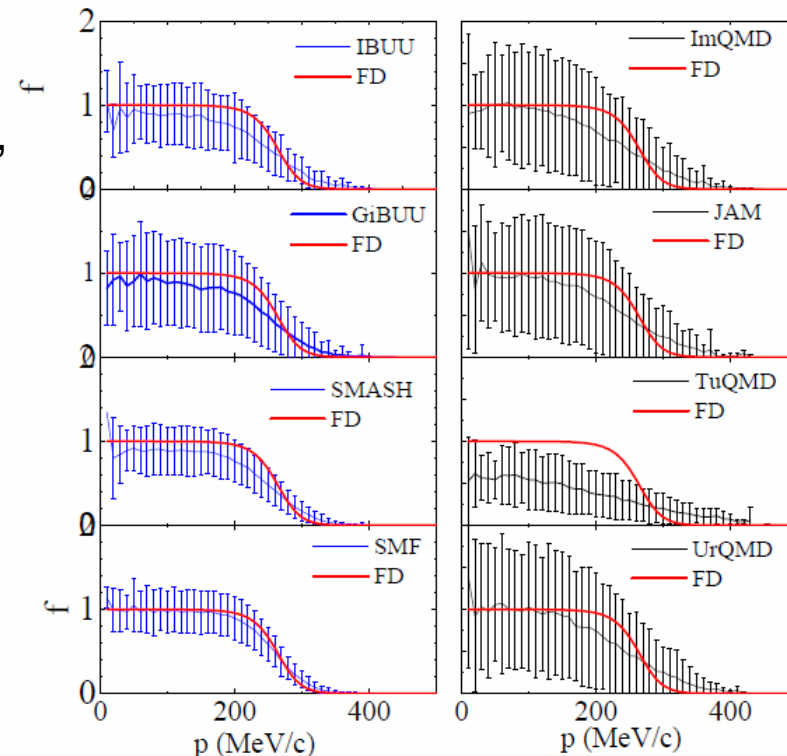
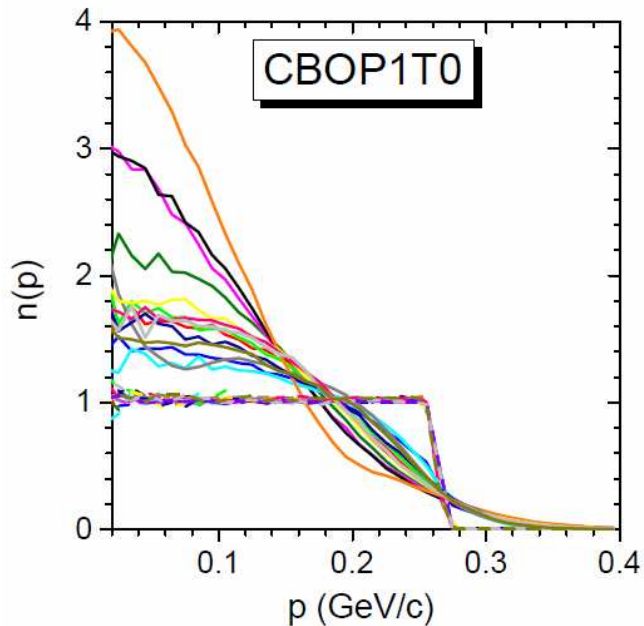
occupation probability $f(r,p,t)$

local

- but realistically averaged over a volume
- often very large, non-localities
- fluctuations!

consequence: evolution to a MB distribution,
 $f(p) > 1$

prescription: $f \leq 1$



how much this affects a transport simulation not clear,
 very likely in the initial stages, e.g. pre-eq emission

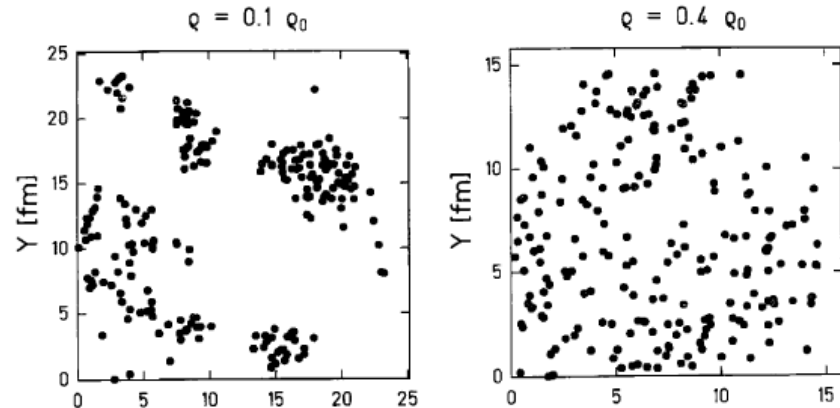
Fluctuations: biggest differences between families of codes and implementation of codes

important: **yes!**

indirect: blocking, mf propagation

direct: fragments formation

test also
fluctuations
and
fragmentation



how treated:

BV-like \rightarrow Boltzmann-Langevin eq.
realizations: BOB, SMF, BLOB

MD-like: damped classical fluctuations
parameter Dx of wave packets

light clusters: another problem, \rightarrow tomorrow afternoon.

freeze out:

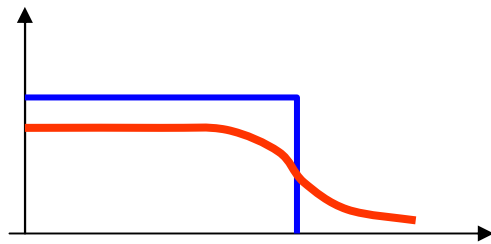
assumption of a completely equilibrated primary fragment is probably too naive

there is still collective motion: expansion

perhaps a differential freeze-out,
surface layer of an expanding source

→ see e.g. Natowitz experiments
check with transport models

short range correlations:



proposed treatments:

1. initialize momentum distribution
 - but has to active at every moment
2. calculate correlation energy in nuclear matter and use this as a part of the potential energy
 - does not generate high energy particles
3. three body collisions, to conserve energy
 - difficult

